

MC@NLO for the hadronic decay of Higgs bosons in associated production with vector bosons

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MC@NLO for the hadronic decay of Higgs bosons in associated production with vector bosons

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ABSTRACT: In this article we describe simulations of the hadronic decay of Higgs bosons produced in association with vector bosons at linear and hadronic colliders. We use the Monte Carlo at next-to-leading-order (MC@NLO) matching prescription with the Herwig++ event generator to predict various spectra of the resulting $b\bar{b}$ pairs and compare our results with leading order and matrix element correction predictions.

KEYWORDS: QCD Phenomenology

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1 Introduction

The Higgs boson is an elusive particle which couples to particles according to their mass. The dominant production mechanisms of Higgs bosons at hadron colliders is from gluon-gluon fusion and vector boson fusion. In this paper, we consider another mechanism which may be more relevant to an experimental search. This is the associated production of Higgs bosons with vector bosons. This production mechanism, $p\bar{p} \rightarrow WH/ZH + X$, is the most promising discovery channel for a light Standard Model (SM) Higgs boson at the Tevatron. This is because the Higgs, which decays predominantly into $b\bar{b}$ pairs, can be tagged by the associated vector boson.

Such processes can be simulated in parton shower generators which resum soft and collinear leading logarithmic, as well as an important subset of next-to-leading logarithmic contributions to all orders. These simulations can further be improved by matching the parton shower to higher order matrix elements. One way in which this is done in the generic parton shower generators is through the use of the matrix element correction [1] which generates harder emissions in regions outside the reach of the parton shower at a

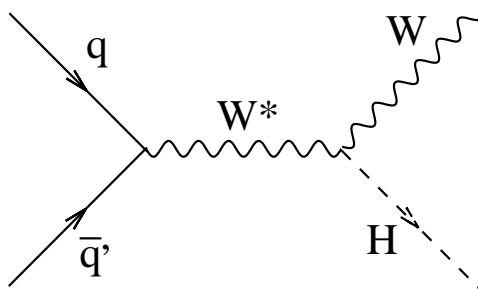


Figure 1. Associated WH boson production.

rate given by the matrix element. A more rigorous matching procedure is the `MC@NLO` method [2–6] which has been implemented for a multitude of processes in the `HERWIG` event generator [7] and for some processes [8, 9] in its successor the `Herwig++` event generator [10, 11]. More recently, another matching method was proposed, called the `POWHEG` (Positive Weighted Hardest Emission Generator) [12, 13], which achieves the same aim as `MC@NLO`, with the creation of positive weighted events and is furthermore independent of the shower generator used. The `POWHEG` method has been applied to Z pair hadroproduction [14], heavy flavour production [15], e^+e^- annihilation to hadrons [16], Drell-Yan vector boson production [17, 18], top pair production at the ILC [19], heavy vector boson production [9] and most recently Higgs boson production [20, 21].

In this paper, we aim to simulate the NLO hadronic decay of the light Higgs boson produced in association with a vector boson using the `MC@NLO` method. The parton shower generator we will be employing is `Herwig++`. In section 2, we first discuss the cross-sections for associated production of the Higgs boson with a vector boson and its subsequent hadronic decay rate. We then discuss the application of the `MC@NLO` method to the decay in section 3. In section 4, we show some comparative distributions obtained from the parton shower and in section 5 we summarize our conclusions. Finally, it should be noted that in this paper, we do not apply the `MC@NLO` method to the initial state emissions.

2 Cross-sections and decay rates

2.1 Associated Higgs production with a W boson from $q\bar{q}$ annihilation

The process $q(p_q) + \bar{q}'(p_{\bar{q}'}) \rightarrow W^* \rightarrow W(p_W) + H(p_H)$ is illustrated in figure 1. If we define the center of mass energy squared of the partonic system by

$$s = (p_q + p_{\bar{q}'})^2, \tag{2.1}$$

we have for the differential cross-section [22],

$$\frac{d\sigma_p}{d\cos\theta^*} = \frac{G_F^2 M_W^2 V_{q\bar{q}'}^2 \beta_W \gamma_W M_W^3}{\cos^2\theta_W 48\pi s^{3/2}} \left(\frac{s + M_W^2 - M_H^2}{s - M_W^2} \right)^2 [2(1 - \beta_W^2) + \beta_W^2 \sin^2\theta^*], \tag{2.2}$$

where θ^* is the angle between the W boson and the quark in the partonic CMF, G_F is the Fermi coupling constant, θ_W is the Weinberg mixing angle, M_W and M_H are respectively

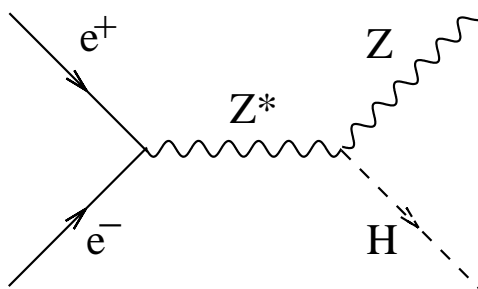


Figure 2. Associated ZH production.

the W and Higgs boson masses and $V_{q\bar{q}'}$ is the relevant CKM matrix element. β_W is the speed of the W boson in the partonic CMF and is given by

$$\beta_W = \frac{\sqrt{[s - (M_W + M_H)^2][s - (M_W - M_H)^2]}}{s - M_H^2 + M_W^2}. \quad (2.3)$$

The relativistic boost factor γ_W is $(1 - \beta_W^2)^{-1/2}$. Equation 2.2 when integrated gives for the total partonic cross-section,

$$\sigma_p = \frac{G_F^2 M_W^2 V_{q\bar{q}'}^2 \beta_W \gamma_W M_W^3}{\cos^2 \theta_W 12\pi s^{3/2}} \left(\frac{s + M_W^2 - M_H^2}{s - M_W^2} \right)^2 \left[1 - \frac{2\beta_W^2}{3} \right]. \quad (2.4)$$

Convolving this with parton distribution functions (PDFs), we obtain the hadronic cross-section as

$$\sigma_h = \int_{\tau}^1 dx_1 \int_{\tau/x_1}^1 dx_2 [f_q(x_1, M_W^2) f_{\bar{q}'}(x_2, M_W^2) + x_1 \leftrightarrow x_2] \sigma_p, \quad (2.5)$$

where x_1, x_2 are the momentum fractions of the incoming partons and taking S as the hadronic beam-beam center-of-mass energy, we have $s = x_1 x_2 S$ and $\tau = \frac{M_W^2}{S}$.

2.2 Associated Higgs production with a Z boson from e^+e^- annihilation

The differential cross-section for the process $e^+e^- \rightarrow Z^* \rightarrow ZH$ is given by [23]

$$\begin{aligned} \frac{d\sigma}{d\cos\theta^*} &= \frac{G_F^2 M_W^2 \beta_Z \gamma_Z M_Z^3}{\cos^2 \theta_W 32\pi s^{3/2}} \left(\frac{s + M_Z^2 - M_H^2}{s - M_Z^2} \right)^2 \times \\ &\times [1 - 4\sin^2 \theta_W + 8\sin^4 \theta_W] [2(1 - \beta_Z^2) + \beta_Z^2 \sin^2 \theta^*], \end{aligned} \quad (2.6)$$

where θ^* is the angle between the electron and the Z boson in the CMF, s is the center-of-mass energy and γ_Z and β_Z are obtained from the analogous expressions for W production in section 2.1 by substituting M_Z for M_W . Equation 2.6 integrated over $\cos\theta^*$ gives

$$\begin{aligned} \sigma(e^+e^- \rightarrow ZH) &= \frac{G_F^2 M_W^2 \beta_Z \gamma_Z M_Z^3}{\cos^2 \theta_W 8\pi s^{3/2}} \left(\frac{s + M_Z^2 - M_H^2}{s - M_Z^2} \right)^2 \times \\ &\times [1 - 4\sin^2 \theta_W + 8\sin^4 \theta_W] \left[1 - \frac{2}{3}\beta_Z^2 \right]. \end{aligned} \quad (2.7)$$

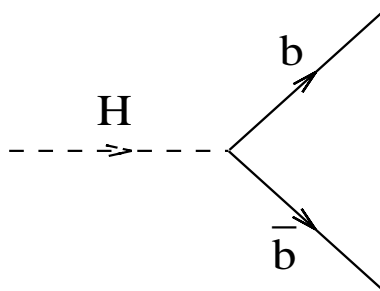


Figure 3. Lowest order Higgs hadronic decay rate.

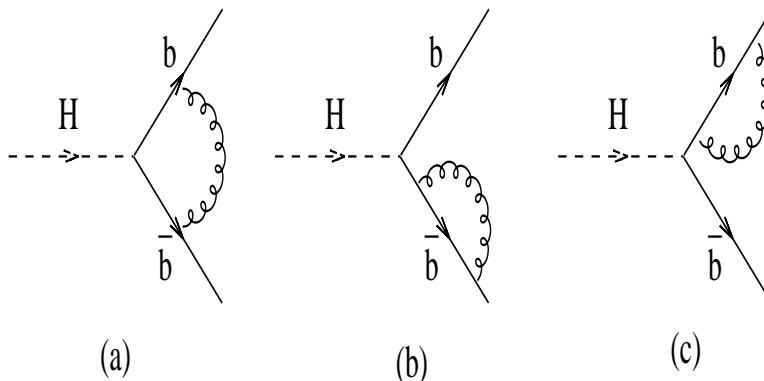


Figure 4. Virtual corrections (a): vertex correction (b), (c): self-energy corrections.

2.3 Higgs boson decay to $b\bar{b}$ pairs

2.3.1 Lowest order decay rate

The lowest-order decay rate for this process, illustrated in figure 3, is given by

$$\Gamma_B(H \rightarrow b\bar{b}) = \frac{3G_F m_b^2}{4\sqrt{2}\pi} M_H \beta_0^3, \quad (2.8)$$

where m_b is the mass of the bottom quark and $\beta_0 = \sqrt{1 - \frac{4m_b^2}{M_H^2}}$.

2.3.2 Virtual radiative corrections

If one uses the on-shell renormalization scheme, the self-energy diagrams in figures 4(b) and 4(c) are cancelled by counter-term diagrams leaving us with the vertex correction in figure 4(a) and its counter-term. When evaluated in the massive gluon regularization scheme, the final result is [24],

$$\begin{aligned} \Gamma_V = \frac{\alpha_S C_F}{\pi} \Gamma_B \left[- \left(\frac{1 + \beta_0^2}{2\beta_0} \ln \frac{1 + \beta_0}{1 - \beta_0} \right) \ln \frac{m_b^2}{\mu^2} \right. \\ \left. + \frac{1 + \beta_0^2}{\beta_0} \left\{ \text{Li}_2 \left(\frac{1 - \beta_0}{1 + \beta_0} \right) + \ln \frac{1 + \beta_0}{2\beta_0} \ln \frac{1 + \beta_0}{1 - \beta_0} - \frac{1}{4} \ln^2 \frac{1 + \beta_0}{1 - \beta_0} + \frac{\pi^2}{3} \right\} \right. \\ \left. + \frac{1 - \beta_0^2}{\beta_0} \ln \frac{1 + \beta_0}{1 - \beta_0} - 1 \right], \quad (2.9) \end{aligned}$$

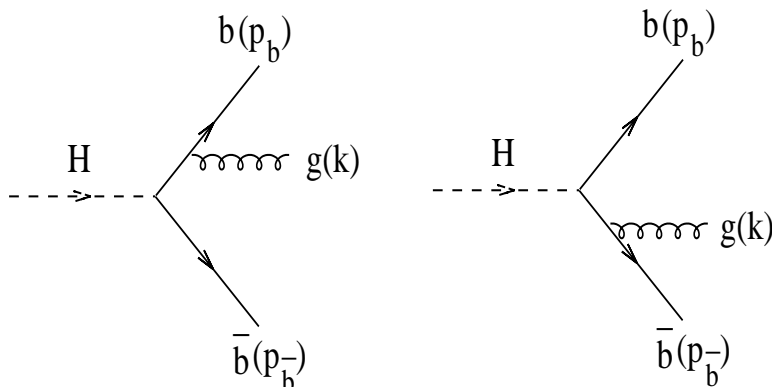


Figure 5. Real gluon emission.

where we have introduced the gluon mass μ to regulate the infrared singularity in figure 4(a).

2.3.3 Real emission corrections

Using the following relations for the energy fractions of the b and \bar{b} quarks in terms of the parton momenta in the Higgs rest frame,

$$\begin{aligned} x_b &= 1 - \frac{2p_{\bar{b}} \cdot k}{M_H^2}, \\ x_{\bar{b}} &= 1 - \frac{2p_b \cdot k}{M_H^2}, \end{aligned} \quad (2.10)$$

we have for the real emission decay rate, the following expression:

$$\begin{aligned} \frac{d\Gamma_R}{dx_b dx_{\bar{b}}} &= \Gamma_B \frac{\alpha_S C_F}{2\pi} \mathcal{M} \\ &= \Gamma_B \frac{\alpha_S C_F}{2\pi \beta_0^3} \left[\frac{(1-x_b)^2 + (1-x_{\bar{b}})^2 + 2(1-2\rho)(x_b+x_{\bar{b}}-1-4\rho)}{(1-x_b)(1-x_{\bar{b}})} \right. \\ &\quad \left. + 4\rho \left\{ \frac{1}{1-x_b} + \frac{1}{1-x_{\bar{b}}} \right\} - 2\rho(1-4\rho) \left\{ \frac{1}{(1-x_b)^2} + \frac{1}{(1-x_{\bar{b}})^2} \right\} + 2 \right], \end{aligned} \quad (2.11)$$

where $\rho = \frac{m_b^2}{M_H^2}$. This can be integrated in the massive gluon scheme to get,

$$\begin{aligned} \Gamma_R &= \Gamma_B \frac{\alpha_S C_F}{2\pi} \left[\left(\frac{1+\beta_0^2}{2\beta_0} \ln \frac{1+\beta_0}{1-\beta_0} \right) \ln \frac{m_b^2}{\mu^2} \right. \\ &\quad + \frac{1+\beta_0^2}{2\beta_0} \left\{ 3\text{Li}_2 \left(\frac{1-\beta_0}{1+\beta_0} \right) + 2\text{Li}_2 \left(-\frac{1-\beta_0}{1+\beta_0} \right) - \ln \frac{2}{1+\beta_0} \ln \frac{1+\beta_0}{1-\beta_0} \right. \\ &\quad \left. + \frac{1}{4} \ln^2 \frac{1+\beta_0}{1-\beta_0} + \ln \frac{1+\beta_0}{2\beta_0} \ln \frac{1+\beta_0}{1-\beta_0} - \frac{\pi^2}{3} \right\} - 3 \ln \frac{4}{1-\beta_0^2} - 4 \ln \beta_0 \\ &\quad \left. + \frac{1}{16\beta_0^3} \{ 3 + 2\beta_0^2 + 3\beta_0^4 \} \ln \frac{1+\beta_0}{1-\beta_0} + \frac{1}{8\beta_0^2} \{-3 + 29\beta_0^2\} \right]. \end{aligned} \quad (2.12)$$

Summing this with Γ_B and Γ_V in equation 2.9, the dependence of the gluon mass μ disappears to give,

$$\Gamma_{\text{NLO}} = \Gamma_B \left[1 + \frac{\alpha_S C_F}{\pi} \Pi \right], \quad (2.13)$$

where

$$\begin{aligned} \Pi = \frac{1}{\beta_0} \left[(1 + \beta_0^2) \left\{ 4\text{Li}_2 \left(\frac{1 - \beta_0}{1 + \beta_0} \right) + 2\text{Li}_2 \left(-\frac{1 - \beta_0}{1 + \beta_0} \right) - 3 \ln \frac{2}{1 + \beta_0} \ln \frac{1 + \beta_0}{1 - \beta_0} \right. \right. \\ \left. \left. - 2 \ln \beta_0 \ln \frac{1 + \beta_0}{1 - \beta_0} \right\} - 3\beta_0 \ln \frac{4}{1 - \beta_0^2} - 4\beta_0 \ln \beta_0 \right] \\ + \frac{1}{16\beta_0^3} [3 + 34\beta_0^2 - 13\beta_0^4] \ln \frac{1 + \beta_0}{1 - \beta_0} + \frac{3}{8\beta_0^2} [-1 + 7\beta_0^2]. \end{aligned} \quad (2.14)$$

Also from equations 2.9, 2.11 and 2.12, we note that Γ_V can be written in terms of the real emission matrix element squared \mathcal{M} as

$$\Gamma_V = \Gamma_B \frac{\alpha_S C_F}{2\pi} \left[\left\{ - \int dx_b dx_{\bar{b}} \mathcal{M} \right\} + 2\Pi_V \right], \quad (2.15)$$

where

$$\begin{aligned} \Pi_V = \frac{1 - \beta_0^2}{\beta_0} \ln \frac{1 + \beta_0}{1 - \beta_0} + \frac{1 + \beta_0^2}{\beta_0} \left\{ 4\text{Li}_2 \left(\frac{1 - \beta_0}{1 + \beta_0} \right) + 2\text{Li}_2 \left(-\frac{1 - \beta_0}{1 + \beta_0} \right) \right. \\ \left. + 2 \ln \frac{1 + \beta_0}{2\beta_0} \ln \frac{1 + \beta_0}{1 - \beta_0} - \ln \frac{2}{1 + \beta_0} \ln \frac{1 + \beta_0}{1 - \beta_0} \right\} - 3 \ln \frac{4}{1 - \beta_0^2} - 4 \ln \beta_0 \\ + \frac{1}{16\beta_0^3} \{3 + 18\beta_0^2 + 3\beta_0^4\} \ln \frac{1 + \beta_0}{1 - \beta_0} + \frac{3}{8\beta_0^2} \{-1 + 7\beta_0^2\}. \end{aligned} \quad (2.16)$$

3 MC@NLO method

The phase space for gluon emission is given in terms of the Dalitz plot variables $x_b, x_{\bar{b}}$ by

$$\lambda(x_b^2 - \rho, x_{\bar{b}}^2 - \rho, (2 - x_b - x_{\bar{b}})^2) \leq 0, \quad (3.1)$$

where the function λ is defined by,

$$\lambda(x, y, z) = x^2 + y^2 + z^2 - 2xy - 2yz - 2xz. \quad (3.2)$$

This is equivalent to the condition

$$(1 - x_b)(1 - x_{\bar{b}})(x_b + x_{\bar{b}} - 1) > \rho(2 - x_b - x_{\bar{b}})^2, \quad (3.3)$$

figure 6 shows the corresponding phase space region where we have labeled as **P** the emission region and **O** the region outside the **P** but in the half-triangle $(1 - x_b)(1 - x_{\bar{b}})(x_b + x_{\bar{b}} - 1) > 0$.

Now using equation 2.11 and 2.15, we can re-write Γ_{NLO} in integral form as,

$$\begin{aligned} \Gamma_{\text{NLO}} = \Gamma_B \int_{\mathbf{P}} dx_b dx_{\bar{b}} \left[2 - \frac{\alpha_S C_F}{2\pi} \{ \mathcal{M} - 4\Pi_V \} + \frac{\alpha_S C_F}{2\pi} \mathcal{M} \right] \\ + \Gamma_B \int_{\mathbf{O}} dx_b dx_{\bar{b}} \left[2 + \frac{\alpha_S C_F}{2\pi} 4\Pi_V \right]. \end{aligned} \quad (3.4)$$

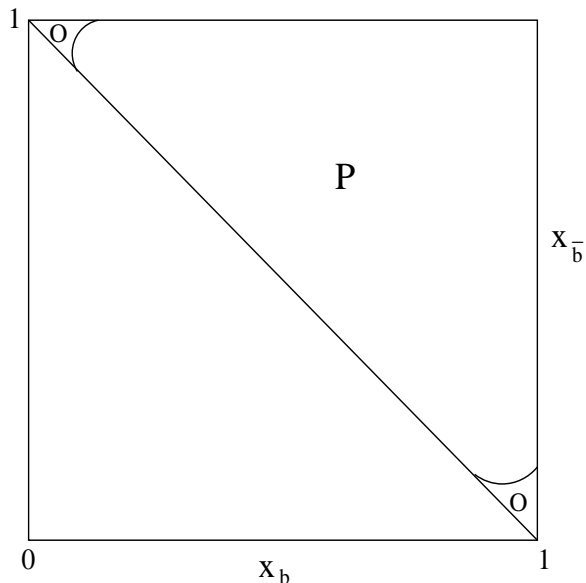


Figure 6. Phase space for gluon emission.

Now, if we define a functional \mathcal{F}_i which represents hadronic final states generated by a parton shower starting from a configuration i , we can write down an overall generating functional for hadrons from Higgs boson decay as

$$\begin{aligned} \mathcal{F} = & \Gamma_B \int_{\mathbf{P}} dx_b dx_{\bar{b}} \left[\mathcal{F}_{b\bar{b}} \left\{ 2 - \frac{\alpha_S C_F}{2\pi} (\mathcal{M} - 4\Pi_V) \right\} + \mathcal{F}_{b\bar{b}g} \frac{\alpha_S C_F}{2\pi} \mathcal{M} \right] \\ & + \Gamma_B \int_{\mathbf{O}} dx_b dx_{\bar{b}} \mathcal{F}_{b\bar{b}} \left[2 + \frac{\alpha_S C_F}{2\pi} 4\Pi_V \right], \end{aligned} \quad (3.5)$$

where $\mathcal{F}_{b\bar{b}}$ is the functional representing shower final states resulting from the process $H \rightarrow b\bar{b}$ and $\mathcal{F}_{b\bar{b}g}$ represents final states from $H \rightarrow b\bar{b}g$.

There are two problems with this functional as it is written above. The first is the highly inefficient sampling that will be required to generate starting $b\bar{b}$ and $b\bar{b}g$ configurations according to the \mathcal{M} since it is divergent in the soft and collinear regions of phase space. The second problem arises because when interfaced with the parton shower, leading order configurations starting with $b\bar{b}$ would radiate quasi-collinear gluons with a distribution given by the parton shower approximation to \mathcal{M} which we shall call \mathcal{M}_C . These are already included in the starting $b\bar{b}g$ configurations. Likewise, some of the $b\bar{b}g$ configurations would include $b\bar{b}$ -like configurations if the gluon is quasi-collinear to either the quark or anti-quark. This problem is often referred to as *double-counting*.

Before we discuss how to solve these problems, let us investigate $\frac{\alpha_S C_F}{2\pi} \mathcal{M}_C$. In the parton shower generator Herwig++, this is the massive quasi-collinear splitting function for the emission of a gluon from a quark [25]:

$$\frac{\alpha_S C_F}{2\pi} \mathcal{M}_C d\tilde{q}^2 dz = \frac{\alpha_S C_F}{2\pi} \frac{d\tilde{q}^2}{\tilde{q}^2} \frac{dz}{1-z} \left[1 + z^2 - \frac{2m_b^2}{z\tilde{q}^2} \right]. \quad (3.6)$$

Here z and \tilde{q} are Herwig++ evolution variables and are respectively the light-cone momentum fraction of the quark after the emission of the gluon and an angular variable related

to the relative transverse momentum of the quark after the emission [26]. The splitting function in equation 3.6 for the b quark can be re-written in terms of the Dalitz plot variables as

$$\frac{\alpha_S C_F}{2\pi} \mathcal{M}_C dx_b dx_{\bar{b}} = \frac{\alpha_S C_F}{2\pi} \frac{dx_b dx_{\bar{b}}}{(1-x_{\bar{b}})\sqrt{x_b^2-4\rho}} \left[\frac{1+z^2}{1-z} - \frac{2\rho}{1-x_{\bar{b}}} \right], \quad (3.7)$$

where if r is defined as

$$r = \frac{1}{2} \left(1 + \frac{\rho}{1+\rho-x_{\bar{b}}} \right), \quad (3.8)$$

then z is given by

$$z = r + \frac{x_b - (2-x_{\bar{b}})r}{\sqrt{x_b^2-4\rho}}. \quad (3.9)$$

Note that $x_b, x_{\bar{b}}$ are given in terms of z, r and \tilde{q} by

$$\begin{aligned} x_b &= 1 - z(1-z) \frac{\tilde{q}^2}{M_H^2}, \\ x_{\bar{b}} &= (2-x_b)r + (z-r)\sqrt{x_b^2-4\rho}. \end{aligned} \quad (3.10)$$

For emission from \bar{b} anti-quark interchange x_b and $x_{\bar{b}}$ in the equations above.

In Herwig++, the quasi-collinear region of phase space covered by the parton shower is defined by imposing the condition,

$$\tilde{q}^2 < \frac{M_H^2}{2} (1 + \sqrt{1-4\rho}). \quad (3.11)$$

The corresponding regions are shown in figure 7 labeled $\mathbf{J}_b, \mathbf{J}_{\bar{b}}$ whilst the unpopulated dead region is labeled \mathbf{D} . Note that region \mathbf{P} in figure 6 corresponds to the union of regions $\mathbf{J}_b, \mathbf{J}_{\bar{b}}$ and \mathbf{D} . Now going back to the functional defined in equation 3.5, if in the region $\mathbf{J} = \mathbf{J}_b \cup \mathbf{J}_{\bar{b}}$, we subtract the parton shower approximation \mathcal{M}_C from the second term in equation 3.5 and add it to the first term we get

$$\begin{aligned} \mathcal{F} &= \Gamma_B \int_{\mathbf{J}} dx_b dx_{\bar{b}} \left[\mathcal{F}_{b\bar{b}} \left\{ 2 - \frac{\alpha_S C_F}{2\pi} (\mathcal{M} - \mathcal{M}_C - 4\Pi_V) \right\} + \mathcal{F}_{b\bar{b}g} \frac{\alpha_S C_F}{2\pi} \{ \mathcal{M} - \mathcal{M}_C \} \right] \\ &+ \Gamma_B \int_{\mathbf{D}} dx_b dx_{\bar{b}} \left[\mathcal{F}_{b\bar{b}} \left\{ 2 - \frac{\alpha_S C_F}{2\pi} (\mathcal{M} - 4\Pi_V) \right\} + \mathcal{F}_{b\bar{b}g} \frac{\alpha_S C_F}{2\pi} \mathcal{M} \right] \\ &+ \Gamma_B \int_{\mathbf{O}} dx_b dx_{\bar{b}} \mathcal{F}_{b\bar{b}} \left[2 + \frac{\alpha_S C_F}{2\pi} 4\Pi_V \right]. \end{aligned} \quad (3.12)$$

This is the MC@NLO method which solves the problem of double-counting in the parton shower regions \mathbf{J} . It also solves the problem of the sampling inefficiency since $\mathcal{M} \rightarrow \mathcal{M}_C$ in the divergent regions and therefore $\mathcal{M} - \mathcal{M}_C$ tends to 0 there. In appendix A, we describe the algorithm used for the evaluation of the above integrals and the generation of events. We also discuss how we regularize some residual divergences by the use of mappings in appendix B.

The procedure followed for event generation for associated Higgs production is outlined below.

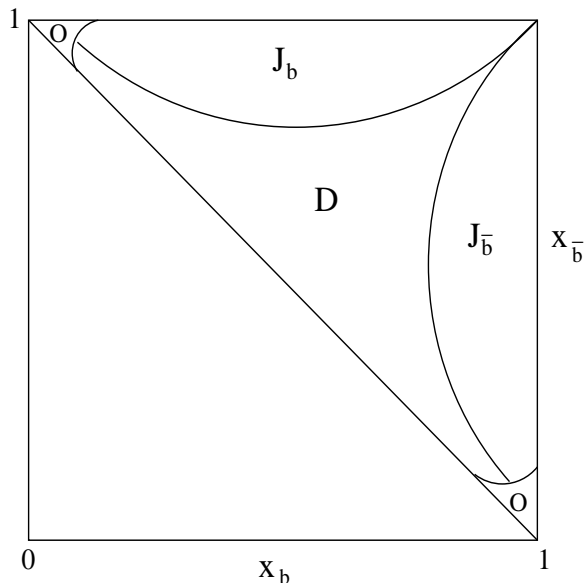


Figure 7. Phase space showing hard **D** and soft/collinear **J_b**, **J_{b̄}** gluon emission regions.

1. For $p\bar{p}$ annihilation use equation 2.5 to distribute the Mandelstam variables x_1, x_2 and the angle θ^* according to the differential cross-section. From these variables, reconstruct the W and Higgs boson four-momenta. For e^+e^- annihilation, use equation 2.6 to distribute the angle θ^* and reconstruct the Z and Higgs boson four-momenta.
2. In the rest frame of the Higgs boson, generate the Dalitz plot variables $x_b, x_{\bar{b}}$ and event weight as described in appendix A. Reconstruct the four-momenta of the quark, anti-quark and gluon in this frame.
3. Boost the four-momenta back to the lab frame.

More details about the MC@NLO method can be found in [2].

4 Results

Following the prescription above, events for associated Higgs boson production of mass 114 GeV and decaying into $b\bar{b}$ pairs with $m_b = 5$ GeV, were generated and interfaced with Herwig++ 2.3.0 [11]. The following approximations were considered:

1. The Herwig++ parton shower interfaced to leading order events (LO),
2. the parton shower interfaced to leading order events which are supplemented by matrix element corrected events in the dead region (ME),
3. the parton shower interfaced to events generated by the MC@NLO method (MC@NLO).

The two processes considered were:

1. associated WH production from $q\bar{q}'$ annihilation at the Tevatron (1.96 TeV),
2. associated ZH production from e^+e^- annihilation at the ILC (0.5 TeV),

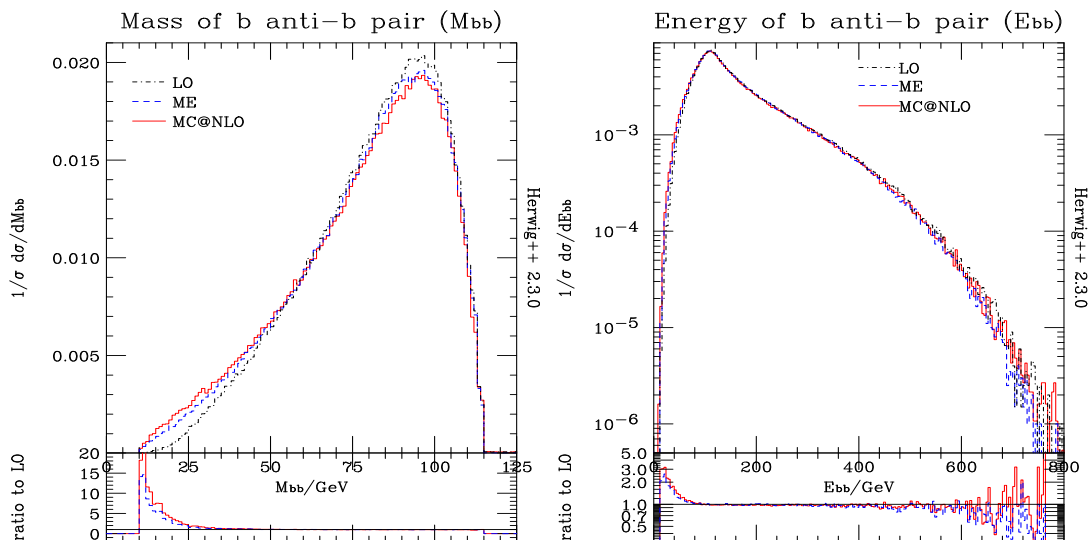


Figure 8. Mass and energy of the $b\bar{b}$ pair ($q\bar{q}'$ annihilation).

with the factorization and renormalization scales set to the Higgs boson mass of 114 GeV.

The following distributions were considered and shown in figures 8–10 for $q\bar{q}'$ annihilation and associated W boson production and figures 11–13 for e^+e^- annihilation and associated Z boson production. In the simulations, only the leptonic decays of vector bosons were considered. All distributions have been normalized to unity.

1. The mass of the $b\bar{b}$ pair before hadronization,
2. the energy of the $b\bar{b}$ pair,
3. the transverse momentum of the $b\bar{b}$ pair with respect to the beam axis,
4. the transverse momentum of the $b\bar{b}$ pair with respect to the direction of the vector boson,
5. the longitudinal momentum of the $b\bar{b}$ pair with respect to the beam axis,
6. the rapidity of the $b\bar{b}$ pair.

From the mass and energy reconstruction plots in figures 8 and 11, we see that the extra gluon radiation simulated by the MC@NLO and matrix element correction methods smooth out the respective peaks due to the production of more low mass and energy $b\bar{b}$ pairs. We also observe as expected that the matrix element correction method slightly underestimates the amount of hard gluon radiation when compared to the MC@NLO distributions.

For $q\bar{q}'$ annihilation, the effect of this extra radiation on the transverse momenta plots is diminished by the initial state radiation from the incoming partons as can be seen in figure 9. The effect is more clearly seen in the corresponding plots for e^+e^- annihilation in figure 12 where there is no initial state radiation. At leading order, the $b\bar{b}$ pairs are produced predominantly at right angles to the beam axis (equation 2.6), hence the peak at high p_T in the LO plot for e^+e^- annihilation. The effect of extra gluon emission is to

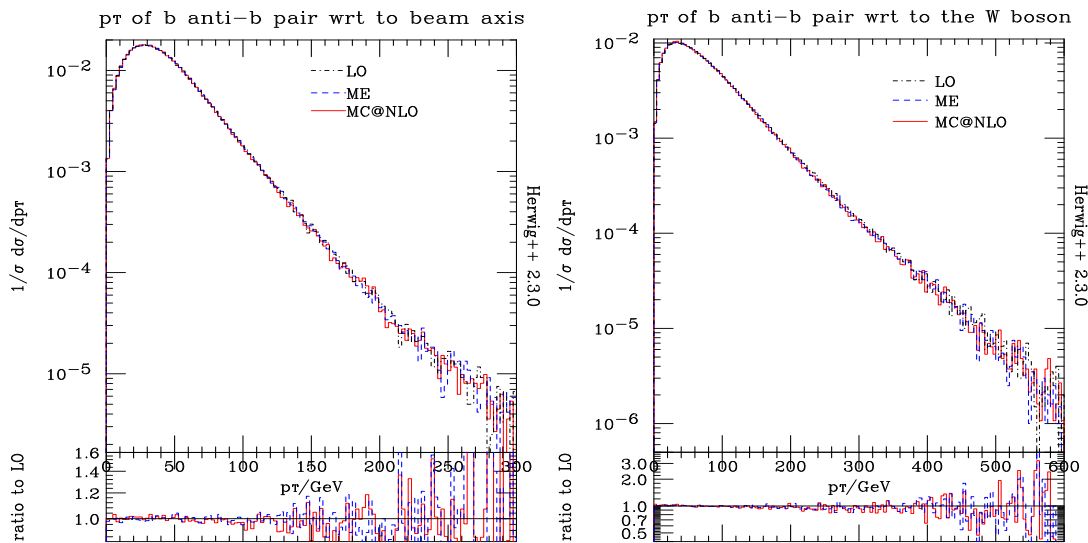


Figure 9. p_T of the $b\bar{b}$ pair w.r.t the beam axis and w.r.t the W boson ($q\bar{q}'$ annihilation).

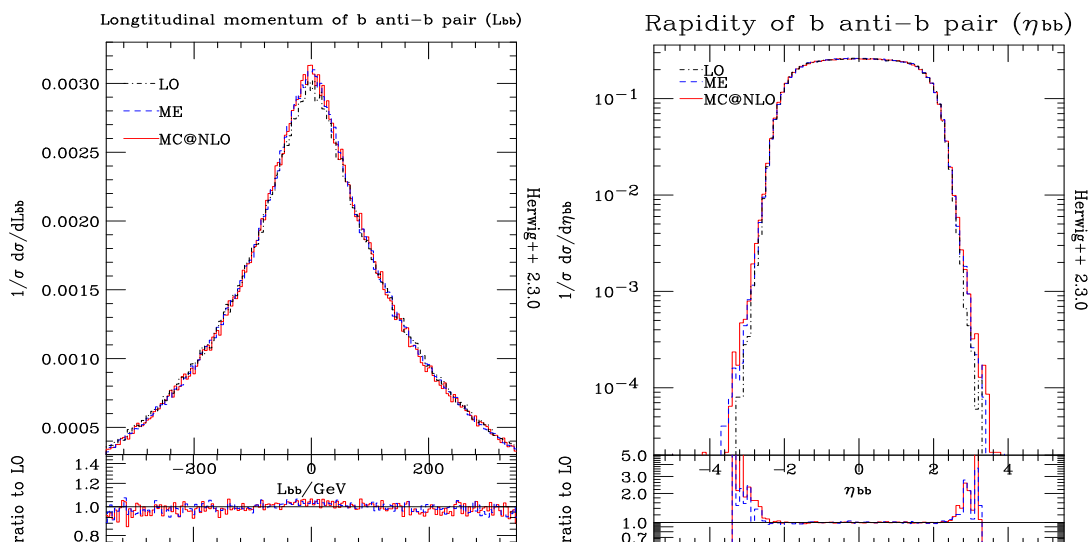


Figure 10. Longitudinal momentum and rapidity of the $b\bar{b}$ pair ($q\bar{q}'$ annihilation).

smooth out the peak as seen in the MC@NLO and ME correction plots. Likewise at leading order, the $b\bar{b}$ pairs are produced back-to-back with the associated Z boson. The effect of gluon radiation is therefore to increase the transverse momentum of the $b\bar{b}$ with respect to the Z boson. This is what is observed in the MC@NLO and ME correction distributions with the latter slightly underestimating the amount of radiation in the tails in comparison with the former.

Finally, in figures 10 and 13, the MC@NLO and ME correction method predict slightly more peaked distributions for the longitudinal momenta around the central value of 0 GeV. This is a result of the loss of energy from the $b\bar{b}$ pairs due to extra gluon radiation. The rapidity distributions show the MC@NLO and ME correction methods predict the production of more high rapidity $b\bar{b}$ pairs in comparison to the leading order prediction. This arises as

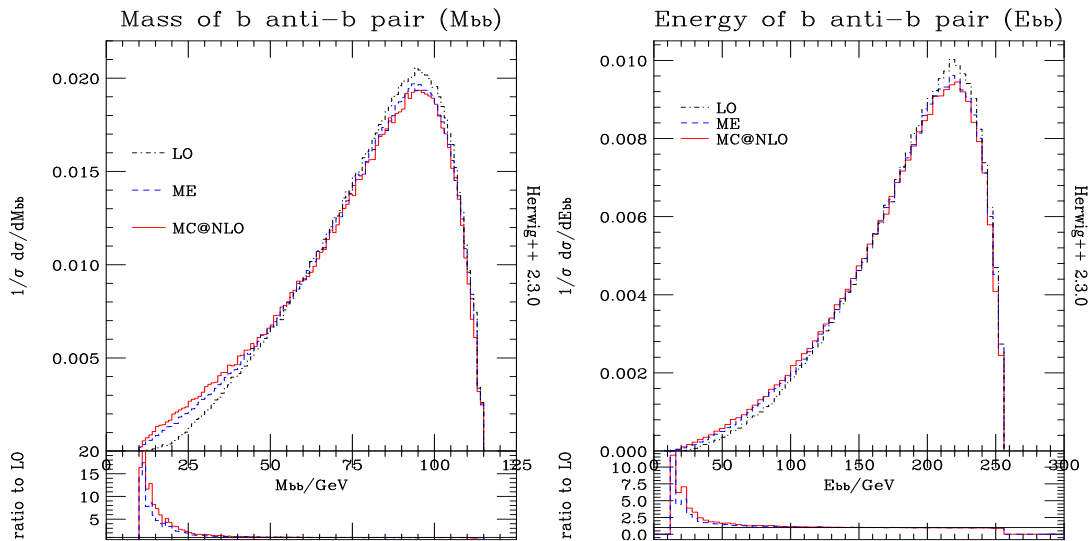


Figure 11. Mass and energy of the $b\bar{b}$ pair (e^+e^- annihilation).

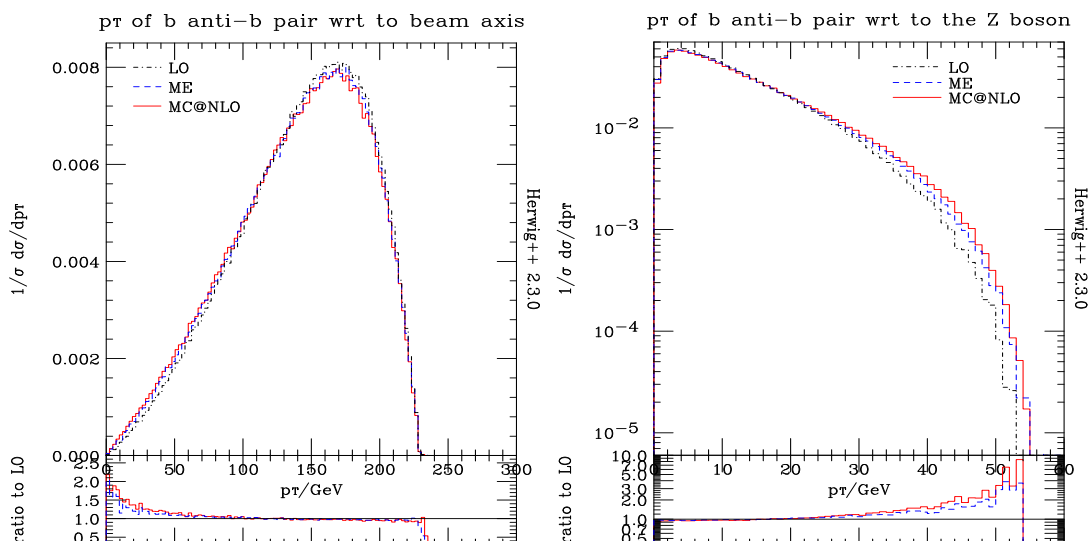


Figure 12. p_T of the $b\bar{b}$ pair w.r.t the beam axis and w.r.t the Z boson (e^+e^- annihilation).

a result of more $b\bar{b}$ pairs being produced at low p_T with respect to the beam axis (figure 12) and therefore higher absolute rapidity. Here the fluctuations of the distributions are such that little or no difference is discernible between the ME correction and MC@NLO plots.

5 Conclusions

In this work we have realized the MC@NLO matching prescription for the decay of Higgs bosons produced in association with vector bosons at both the ILC and hadron colliders. This work was achieved within the framework of the Herwig++ Monte Carlo event generator.

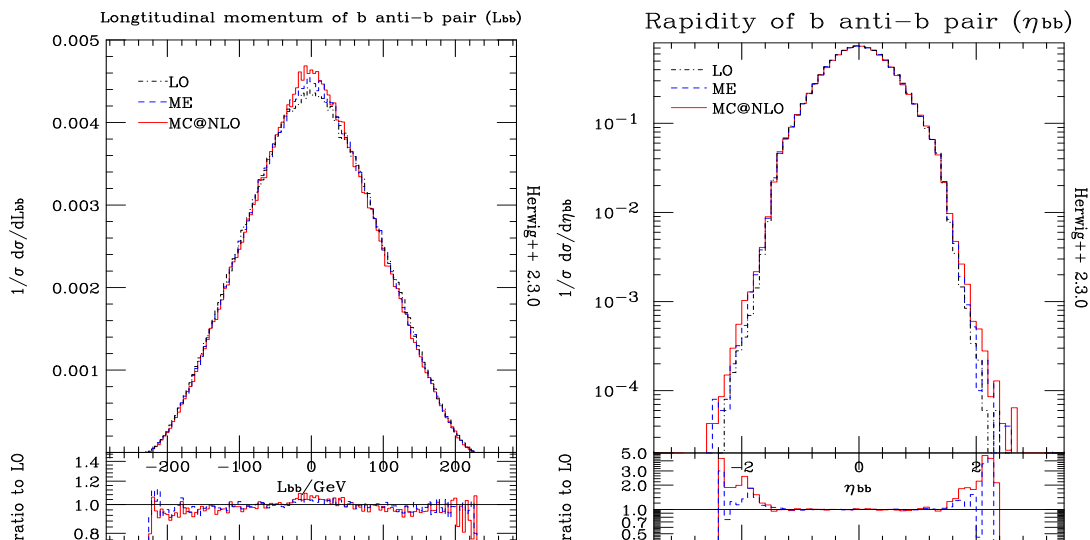


Figure 13. Longitudinal momentum and rapidity of the $b\bar{b}$ pair (e^+e^- annihilation).

We compared the MC@NLO predictions with those obtained via the matrix element correction method as well as leading order predictions. The effects of the hard radiation are visible in the reconstruction plots for the mass and energy of the $b\bar{b}$ pairs resulting from Higgs decay. Also visible is the effect on the longitudinal momenta and rapidity of the $b\bar{b}$ pairs.

Less visible is the effect on the p_T spectra for $q\bar{q}'$ annihilation due to the dominant effects of initial state radiation. For e^+e^- annihilation, the MC@NLO method predicts a softer spectrum for the p_T with respect to the beam axis and a harder spectrum with respect to the associated Z boson.

The algorithm used for these processes will be publicly available with the forthcoming version of Herwig++. Although we have considered the associated production of Higgs and vector bosons in this paper, the MC@NLO algorithm can be interfaced with other Higgs boson production mechanisms, since here we only apply the method to the decay process.

Acknowledgments

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A Monte Carlo algorithm

The integrals in (3.12) can be evaluated using a variety of Monte Carlo methods. In this paper, the ‘Hit or Miss’ Monte Carlo method is used. This is the simplest and oldest form

of Monte Carlo integration and essentially involves finding the area of a region in phase space by integrating over a larger region, a binary function which is 1 in the region and 0 elsewhere. The sampling method used for the points $x_b, x_{\bar{b}}$ is the importance sampling method whereby more samples are taken from regions where the integrand is large and less from regions where it is small. This ensures that the sampled points have the same distribution as the integrand.

The following algorithm summarizes how the starting $b\bar{b}$ and $b\bar{b}g$ configurations were generated according to equation 3.12. In the discussion that follows, we work in the rest frame of a Higgs boson of mass $M_H = 114 \text{ GeV}$ and $\alpha_S = \alpha_S(M_H) = 0.114$.

1. Randomly sample points $x_b, x_{\bar{b}}$, in each of regions $\mathbf{J}_b, \mathbf{J}_{\bar{b}}, \mathbf{D}$ and \mathbf{O} of the phase space and using the ‘Hit Or Miss’ Monte Carlo method, evaluate the 5 integrals, $I_J^{(2)}, I_J^{(3)}, I_D^{(2)}, I_D^{(3)}$ and $I_O^{(2)}$ as well as their absolute sum, I .

$$\begin{aligned}
 I_J^{(2)} &= \int_{\mathbf{J}} dx_b dx_{\bar{b}} \left[2 - \frac{\alpha_S}{2\pi} C_F \{ \mathcal{M} - \mathcal{M}_C - 4\Pi_V \} \right], \\
 I_J^{(3)} &= \int_{\mathbf{J}} dx_b dx_{\bar{b}} \frac{\alpha_S}{2\pi} C_F [\mathcal{M} - \mathcal{M}_C], \\
 I_D^{(2)} &= \int_{\mathbf{D}} dx_b dx_{\bar{b}} \left[2 - \frac{\alpha_S}{2\pi} C_F \{ \mathcal{M} - 4\Pi_V \} \right], \\
 I_D^{(3)} &= \int_{\mathbf{D}} dx_b dx_{\bar{b}} \frac{\alpha_S}{2\pi} C_F \mathcal{M}, \\
 I_O^{(2)} &= \int_{\mathbf{O}} dx_b dx_{\bar{b}} \left[2 + \frac{\alpha_S}{2\pi} C_F 4\Pi_V \right], \\
 I &= | I_J^{(2)} | + | I_J^{(3)} | + | I_D^{(2)} | + | I_D^{(3)} | + | I_O^{(2)} |. \tag{A.1}
 \end{aligned}$$

Note also the maximum values of the integrands in $I_J^{(3)}$ and $I_D^{(3)}$.

2. The eventual proportion of $b\bar{b}$ Monte Carlo events will be determined by the ratio $\frac{|I_J^{(2)}| + |I_D^{(2)}| + |I_O^{(2)}|}{I}$. Likewise, the proportion of $b\bar{b}g$ events in the soft regions $\mathbf{J}_b, \mathbf{J}_{\bar{b}}$ and the hard region \mathbf{D} are determined by the ratios $\frac{|I_J^{(3)}|}{I}$ and $\frac{|I_D^{(3)}|}{I}$ respectively. The algorithm below is then used to importance-sample the $b\bar{b}g$ events so that the corresponding $(x_b, x_{\bar{b}})$ values of the Monte Carlo events have the same distribution as the integrands in $I_J^{(3)}$ and $I_D^{(3)}$:

- (a) For event generation in region \mathbf{L} (where \mathbf{L} is one of \mathbf{D}, \mathbf{J}_b or $\mathbf{J}_{\bar{b}}$), randomly select a point $x_b, x_{\bar{b}}$ in that region.
- (b) Evaluate the absolute value of the integrand in $I_L^{(3)}$ for this point, $|w(x_b, x_{\bar{b}})|$. Is $|w(x_b, x_{\bar{b}})| > R |w_{\max}|$? (R is a random number between 0 and 1 and $|w_{\max}|$ is the maximum value of $|w(x_b, x_{\bar{b}})|$ determined in Step 1).
- (c) If NO, return to (a). If YES, accept the event and set $w^{\text{unw}} = \text{sgn } w(x_b, x_{\bar{b}})$ i.e. $w^{\text{unw}} = 1$ if $w(x_b, x_{\bar{b}})$ is positive and -1 if negative. (In regions \mathbf{J}_b and $\mathbf{J}_{\bar{b}}$, $\mathcal{M} < \mathcal{M}_C$ and hence the integrands and the integral, $I_J^{(3)}$ in these regions are negative). This process is called **unweighting**.

- (d) Repeat the process until the correct proportion of $b\bar{b}$ and $b\bar{b}g$ events have been generated.
- (e) Using the importance-sampled points, obtain an estimate for the integral, $I_L^{(2,3)} = \frac{\sum w^{\text{unw}}}{N} \times I$, where N is the total number of Monte Carlo events generated. We typically use $N = 10^6$.

This is the MC@NLO method. In this way, for a total of N events, the correct proportion of $b\bar{b}$ and $b\bar{b}g$ events with \pm unit weight is generated with the same distribution as the integrands in (A.1). All of these integrals are finite, but the integrands are divergent at isolated points within the integration regions. Before the sampling could be carried out, the divergences in the integrands (which cause problems in the sampling process) had to be taken care of. This is described in section B.

B Divergences and mappings

B.1 Divergences in dead region

In region **D**, the hard matrix element squared \mathcal{M} given in equation 2.11, diverges as $(x_b, x_{\bar{b}}) \rightarrow (1, 1)$. To avoid this divergence, one can map the divergent region into another region in such a way that the divergence is regularized. This is ensured by the fact that the region of integration vanishes as the singularity is approached. There is a double pole in \mathcal{M} at $(x_b, x_{\bar{b}}) = (1, 1)$. To avoid this pole, the region $x_b, x_{\bar{b}} > \frac{3}{4}$ is mapped into a region which includes **D** but whose width vanishes quadratically as $x_b, x_{\bar{b}} \rightarrow 1$ [8, 26]. The mapping is:

$$\begin{aligned} x'_b &= 1 - \left[\frac{1}{4} - (1 - x_b) \right] = \frac{7}{4} - x_b, \\ x'_{\bar{b}} &= 1 - 2(1 - x'_b) \left[\frac{3}{4} - (1 - x_{\bar{b}}) \right] = \frac{5}{8} + \frac{1}{2}x_b + \frac{3}{2}x_{\bar{b}} - 2x_b x_{\bar{b}} \end{aligned} \quad (\text{B.1})$$

when $x_b > x_{\bar{b}} > \frac{3}{4}$. This mapping also introduces an extra weight factor of $2(1 - x'_b)$ in the integrand. Interchange x_b and $x_{\bar{b}}$ in both the mapping and weight factor when $x_{\bar{b}} > x_b > \frac{3}{4}$. Figure 14 shows the region mapped (solid) and the region mapped onto (dashed).

B.2 Divergences in jet regions \mathbf{J}_b and $\mathbf{J}_{\bar{b}}$

In both regions \mathbf{J}_b and $\mathbf{J}_{\bar{b}}$, there is a simple pole in the term $(\mathcal{M} - \mathcal{M}_C)$ at $(x_b, x_{\bar{b}}) = (1, 1)$. In the region $x_b, x_{\bar{b}} > \frac{3}{4}$, a new set of random points are generated which have a weight factor to cancel the divergence. The mapping used in region \mathbf{J}_b where $x_{\bar{b}} > x_b$ is [8]:

$$\begin{aligned} x'_b &= 1 - 0.25r_1, \\ x'_{\bar{b}} &= 1 - (1 - x'_b)r_2 \end{aligned} \quad (\text{B.2})$$

where r_1 and r_2 are random numbers in the range $[0, 1]$. The weight factor for this mapping is $2r_1$. For region $\mathbf{J}_{\bar{b}}$, where $x_b > x_{\bar{b}}$, interchange x_b and $x_{\bar{b}}$ in the mapping. The mapped regions are shown with solid boundaries in figure 15.

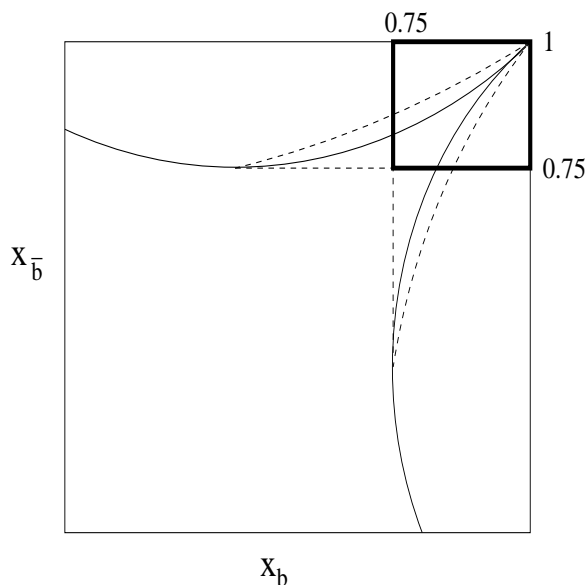


Figure 14. The mapped region (solid) and the region mapped onto (dashed)

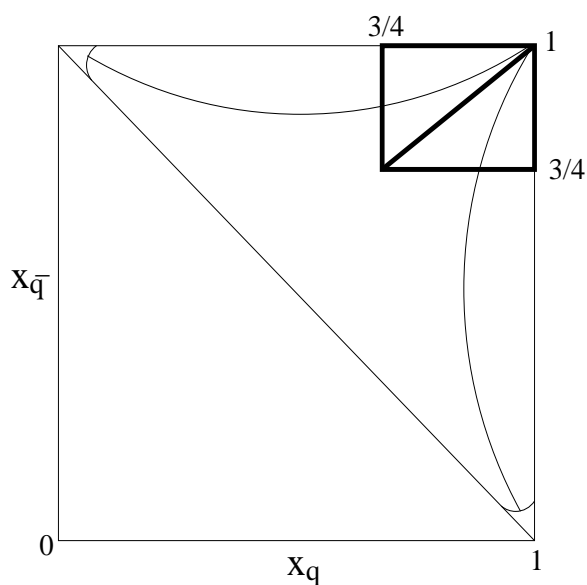


Figure 15. Mapped regions

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